

## Problem 3.40

The most general wave function of a particle in the simple harmonic oscillator potential is

$$\Psi(x, t) = \sum_n c_n \psi_n(x) e^{-iE_n t/\hbar}.$$

Show that the expectation value of position is

$$\langle x \rangle = C \cos(\omega t - \phi),$$

where the real constants  $C$  and  $\phi$  are given by

$$C e^{-i\phi} = \left( \sqrt{\frac{2\hbar}{m\omega}} \right) \sum_{n=0}^{\infty} \sqrt{n+1} c_{n+1}^* c_n.$$

Thus the expectation value of position for a particle in the harmonic oscillator oscillates at the classical frequency  $\omega$  (as you would expect from Ehrenfest's theorem; see problem 3.19(b)). *Hint:* Use Equation 3.114. As an example, find  $C$  and  $\phi$  for the wave function in Problem 2.40.

[Capitalize the “p” to be consistent.]

### Solution

#### Method 1

Use the method of Example 2.5 on page 47 and express the position operator in terms of the promotion and demotion operators,  $\hat{a}_+$  and  $\hat{a}_-$ , respectively.

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}_+ + \hat{a}_-)$$

Representing the  $n$ th eigenstate  $\psi_n(x)$  of the harmonic oscillator as a ket  $|n\rangle$ , the promotion and demotion operators satisfy

$$\hat{a}_+ |n\rangle = \sqrt{n+1} |n+1\rangle$$

$$\hat{a}_- |n\rangle = \sqrt{n} |n-1\rangle.$$

So then

$$\begin{aligned} \langle x \rangle &= \langle \Psi | \hat{x} | \Psi \rangle \\ &= \left( \sum_{q=0}^{\infty} c_q^* e^{iE_q t/\hbar} \langle q | \right) \hat{x} \left( \sum_{n=0}^{\infty} c_n e^{-iE_n t/\hbar} |n\rangle \right) \\ &= \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \langle q | \hat{x} | n \rangle \\ &= \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \langle q | \cdot (\hat{x} | n \rangle). \end{aligned}$$

Use the fact that the eigenstates of the harmonic oscillator are orthonormal:  $\langle q | n \rangle = \delta_{qn}$ .

$$\begin{aligned}
 \langle x \rangle &= \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \langle q | \cdot \left[ \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}_+ + \hat{a}_-) | n \rangle \right] \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \langle q | \cdot (\hat{a}_+ | n \rangle + \hat{a}_- | n \rangle) \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \langle q | \cdot (\sqrt{n+1} | n+1 \rangle + \sqrt{n} | n-1 \rangle) \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} (\sqrt{n+1} \langle q | n+1 \rangle + \sqrt{n} \langle q | n-1 \rangle) \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} (\sqrt{n+1} \delta_{q,n+1} + \sqrt{n} \delta_{q,n-1}) \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} (\sqrt{n+1} \delta_{q,n+1} + \sqrt{n} \delta_{q+1,n}) \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \left[ \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \sqrt{n+1} \delta_{q,n+1} + \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \sqrt{n} \delta_{q+1,n} \right] \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \left[ \sum_{n=0}^{\infty} c_{n+1}^* c_n e^{i(E_{n+1} - E_n)t/\hbar} \sqrt{n+1} + \sum_{q=0}^{\infty} c_q^* c_{q+1} e^{i(E_q - E_{q+1})t/\hbar} \sqrt{q+1} \right]
 \end{aligned}$$

$q$  is just a dummy index and can be replaced with  $n$ .

$$\begin{aligned}
 \langle x \rangle &= \sqrt{\frac{\hbar}{2m\omega}} \left[ \sum_{n=0}^{\infty} c_{n+1}^* c_n e^{i(E_{n+1} - E_n)t/\hbar} \sqrt{n+1} + \sum_{n=0}^{\infty} c_n^* c_{n+1} e^{i(E_n - E_{n+1})t/\hbar} \sqrt{n+1} \right] \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} \left[ c_{n+1}^* c_n e^{i(E_{n+1} - E_n)t/\hbar} + c_n^* c_{n+1} e^{i(E_n - E_{n+1})t/\hbar} \right] \\
 &= 2\sqrt{\frac{\hbar}{2m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} \left\{ \frac{c_{n+1}^* c_n e^{i(E_{n+1} - E_n)t/\hbar} + [c_{n+1}^* c_n e^{i(E_{n+1} - E_n)t/\hbar}]^*}{2} \right\} \\
 &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} \operatorname{Re} c_{n+1}^* c_n e^{i(E_{n+1} - E_n)t/\hbar}
 \end{aligned}$$

Note that for a complex number  $z$ ,  $\operatorname{Re} z = (z + z^*)/2$ . The energy of the  $n$ th eigenstate of the harmonic oscillator is

$$E_n = \left( n + \frac{1}{2} \right) \hbar\omega,$$

so  $E_{n+1} - E_n = \hbar\omega$ .

As a result,

$$\begin{aligned}\langle x \rangle &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} \operatorname{Re} c_{n+1}^* c_n e^{i\omega t} \\ &= \operatorname{Re} e^{i\omega t} \left( \sqrt{\frac{2\hbar}{m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} c_{n+1}^* c_n \right).\end{aligned}$$

Define the complex number in parentheses to be  $Ce^{-i\phi}$ , where  $C$  and  $\phi$  are real numbers.

$$\begin{aligned}\langle x \rangle &= \operatorname{Re} e^{i\omega t} (Ce^{-i\phi}) \\ &= C \operatorname{Re} e^{i(\omega t - \phi)}\end{aligned}$$

Therefore,

$$\langle x \rangle = C \cos(\omega t - \phi).$$

**Method 2**

This is another way to prove the result, using the boxed formula for  $\langle n | \hat{x} | n' \rangle$  in Problem 3.39.

$$\langle n | \hat{x} | n' \rangle = \sqrt{\frac{\hbar}{2m\omega}} \left( \sqrt{n} \delta_{n,n'+1} + \sqrt{n'} \delta_{n',n+1} \right)$$

Consequently,

$$\begin{aligned} \langle x \rangle &= \langle \Psi | \hat{x} | \Psi \rangle \\ &= \langle \Psi | \hat{I} \hat{x} \hat{I} | \Psi \rangle \\ &= \langle \Psi | \left( \sum_{j=0}^{\infty} |j\rangle \langle j| \right) \hat{x} \left( \sum_{k=0}^{\infty} |k\rangle \langle k| \right) | \Psi \rangle \\ &= \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \langle \Psi | j \rangle \langle j | \hat{x} | k \rangle \langle k | \Psi \rangle \\ &= \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \langle \Psi | j \rangle \left[ \sqrt{\frac{\hbar}{2m\omega}} \left( \sqrt{j} \delta_{j,k+1} + \sqrt{k} \delta_{k,j+1} \right) \right] \langle k | \Psi \rangle \\ &= \sqrt{\frac{\hbar}{2m\omega}} \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \langle \Psi | j \rangle \left( \sqrt{j} \delta_{j,k+1} + \sqrt{k} \delta_{k,j+1} \right) \langle k | \Psi \rangle \\ &= \sqrt{\frac{\hbar}{2m\omega}} \left( \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \sqrt{j} \langle \Psi | j \rangle \langle k | \Psi \rangle \delta_{j,k+1} + \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \sqrt{k} \langle \Psi | j \rangle \langle k | \Psi \rangle \delta_{k,j+1} \right) \\ &= \sqrt{\frac{\hbar}{2m\omega}} \left( \sum_{k=0}^{\infty} \sqrt{k+1} \langle \Psi | k+1 \rangle \langle k | \Psi \rangle + \sum_{j=0}^{\infty} \sqrt{j+1} \langle \Psi | j \rangle \langle j+1 | \Psi \rangle \right). \end{aligned}$$

$j$  is just a dummy index and can be replaced with  $k$ .

$$\begin{aligned} &= \sqrt{\frac{\hbar}{2m\omega}} \left( \sum_{k=0}^{\infty} \sqrt{k+1} \langle \Psi | k+1 \rangle \langle k | \Psi \rangle + \sum_{k=0}^{\infty} \sqrt{k+1} \langle \Psi | k \rangle \langle k+1 | \Psi \rangle \right) \\ &= \sqrt{\frac{\hbar}{2m\omega}} \sum_{k=0}^{\infty} \sqrt{k+1} \left( \langle \Psi | k+1 \rangle \langle k | \Psi \rangle + \langle \Psi | k \rangle \langle k+1 | \Psi \rangle \right) \\ &= 2\sqrt{\frac{\hbar}{2m\omega}} \sum_{k=0}^{\infty} \sqrt{k+1} \left[ \frac{\langle \Psi | k+1 \rangle \langle k | \Psi \rangle + (\langle \Psi | k+1 \rangle \langle k | \Psi \rangle)^*}{2} \right] \\ &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{k=0}^{\infty} \sqrt{k+1} \operatorname{Re} \langle \Psi | k+1 \rangle \langle k | \Psi \rangle \end{aligned}$$

Substitute the formula for  $\Psi(x, t)$ .

$$\begin{aligned}
 \langle x \rangle &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{k=0}^{\infty} \sqrt{k+1} \operatorname{Re} \left( \sum_{q=0}^{\infty} c_q^* e^{iE_q t/\hbar} \langle q | \right) |k+1\rangle \langle k| \left( \sum_{n=0}^{\infty} c_n e^{-iE_n t/\hbar} |n\rangle \right) \\
 &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{k=0}^{\infty} \sqrt{k+1} \operatorname{Re} \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \langle q | k+1 \rangle \langle k | n \rangle \\
 &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{k=0}^{\infty} \sqrt{k+1} \operatorname{Re} \sum_{q=0}^{\infty} \sum_{n=0}^{\infty} c_q^* c_n e^{i(E_q - E_n)t/\hbar} \delta_{q, k+1} \delta_{kn} \\
 &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{k=0}^{\infty} \sqrt{k+1} \operatorname{Re} \sum_{n=0}^{\infty} c_{k+1}^* c_n e^{i(E_{k+1} - E_n)t/\hbar} \delta_{kn} \\
 &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} \operatorname{Re} c_{n+1}^* c_n e^{i(E_{n+1} - E_n)t/\hbar}
 \end{aligned}$$

The energy of the  $n$ th eigenstate of the harmonic oscillator is

$$E_n = \left( n + \frac{1}{2} \right) \hbar\omega,$$

so  $E_{n+1} - E_n = \hbar\omega$ .

$$\begin{aligned}
 \langle x \rangle &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} \operatorname{Re} c_{n+1}^* c_n e^{i\omega t} \\
 &= \operatorname{Re} e^{i\omega t} \left( \sqrt{\frac{2\hbar}{m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} c_{n+1}^* c_n \right)
 \end{aligned}$$

Define the complex number in parentheses to be  $C e^{-i\phi}$ , where  $C$  and  $\phi$  are real numbers.

$$\begin{aligned}
 \langle x \rangle &= \operatorname{Re} e^{i\omega t} (C e^{-i\phi}) \\
 &= C \operatorname{Re} e^{i(\omega t - \phi)}
 \end{aligned}$$

Therefore,

$$\langle x \rangle = C \cos(\omega t - \phi).$$

**Problem 2.40**

In Problem 2.40 the wave function is

$$\begin{aligned}\Psi(x, t) &= c_0\psi_0(x)e^{-iE_0t/\hbar} + c_1\psi_1(x)e^{-iE_1t/\hbar} + c_2\psi_2(x)e^{-iE_2t/\hbar} \\ &= \frac{3}{5}\psi_0(x)e^{-i\omega t/2} - \frac{2\sqrt{2}}{5}\psi_1(x)e^{-3i\omega t/2} + \frac{2\sqrt{2}}{5}\psi_2(x)e^{-5i\omega t/2}.\end{aligned}$$

Therefore,

$$\begin{aligned}Ce^{-i\phi} &= \sqrt{\frac{2\hbar}{m\omega}} \sum_{n=0}^{\infty} \sqrt{n+1} c_{n+1}^* c_n \\ &= \sqrt{\frac{2\hbar}{m\omega}} \left( \sqrt{0+1} c_{0+1}^* c_0 + \sqrt{1+1} c_{1+1}^* c_1 + \underbrace{\sqrt{2+1} c_{2+1}^* c_2}_{=0} + \underbrace{\sqrt{3+1} c_{3+1}^* c_3}_{=0} + \dots \right) \\ &= \sqrt{\frac{2\hbar}{m\omega}} \left( c_1^* c_0 + \sqrt{2} c_2^* c_1 \right) \\ &= \sqrt{\frac{2\hbar}{m\omega}} \left[ \left( -\frac{2\sqrt{2}}{5} \right) \left( \frac{3}{5} \right) + \sqrt{2} \left( \frac{2\sqrt{2}}{5} \right) \left( -\frac{2\sqrt{2}}{5} \right) \right] \\ &= \sqrt{\frac{2\hbar}{m\omega}} \left( -\frac{14\sqrt{2}}{25} \right) \\ &= -\frac{28}{25} \sqrt{\frac{\hbar}{m\omega}},\end{aligned}$$

which means

$$C = \frac{28}{25} \sqrt{\frac{\hbar}{m\omega}} \quad \text{and} \quad \phi = \pi + 2\pi n,$$

where  $n$  is any integer.